



The cosine method to Gardner equation and (2+1)- dimensional breaking soliton system

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Received 02 May 2016; accepted 4 February 2017

Abstract

In this letter, we established a traveling wave solution by using cosine function algorithm for Gardner equation and (2+1)-dimensional breaking soliton system. The cosine method is used to obtain the exact solution.

Key words: Gardner equation; cosine function method; exact solution;
(2+1)dimensional breaking soliton system

2010 AMS Mathematics Subject Classification : 65L11; 46E22; 65F25.

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1 introduction

Gardner equation known as the mixed $kdv - mkdv$ equation is very widely studied in various areas of physics that includes plasma physics, Fluid dynamic, Quantum Field Theory, solid state physics and other [3] the Gardner equation is solved by sin-cosine function method [2]. the breaking soliton system was used to describes the (2+1)-dimensional interaction of Riemann wave propagated along the y-axis with long wave propagated along the x- axis and it seems to have been investigated extensively where over lapping solutions have been derived. this system is solved to Generalized jacobii elliptic function method [4] and the (G'/G) -Expansion method [5].

2 The cosine-function method

Consider the nonlinear partial differential equation in the form

$$F(u, u_t, u_x, u_{xx}, u_{xxt}) = 0, \quad (1)$$

Where $u(x, t)$ is the solution of nonlinear partial differential equation Eq.(1). we use the

$$u(x, t) = f(\xi), \quad (2)$$

Where $\xi = x - ct$. This enables us to use the following changes:

$$\begin{aligned} \frac{\partial}{\partial t}(\dots) &= -c \frac{d}{d\xi}, & \frac{\partial}{\partial x}(\dots) &= \frac{d}{d\xi}(\dots), \\ \frac{\partial^2}{\partial x^2}(\dots) &= \frac{d^2}{d\xi^2}(\dots), & \dots & \end{aligned} \quad (3)$$

Using Eq.(3) to transfer the nonlinear partial differential equation Eq.(??) to nonlinear ordinary differential equation

$$G(f, f', f'', f''', \dots) = 0. \quad (4)$$

The solution of Eq.(4) can be expressed in the form:

$$f(\xi) = \lambda \cos^\beta(\mu\xi), \quad |\xi| \leq \frac{\pi}{2\mu}, \quad (5)$$

Where λ , β and μ are unknown parameters which will be determined. Then we have:

$$\begin{aligned} f' &= \frac{df(\xi)}{d\xi} = \lambda\beta\mu \cos^{\beta-1}(\mu\xi) \sin(\mu\xi), \\ f'' &= \frac{d^2f(\xi)}{d\xi^2} = -\lambda\beta\mu^2 \cos^\beta(\mu\xi) + \lambda\mu^2\beta(\beta-1) \cos^{\beta-2}(\mu\xi) \\ &\quad -\lambda\mu^2\beta(\beta-1) \cos^\beta(\mu\xi). \end{aligned} \quad (6)$$

Substituting Eq.(6) into the nonlinear ordinary differential equation Eq.(4) gives a trigonometric of terms. To determine the parameters first balancing the exponents of each pair of cosine to determine α . Then we collect all terms with the same power in $\cos^\beta(\mu\xi)$ and put to zero their coefficients to get a system of algebraic equations among the unknown β , λ and μ . Now, the problem is reduced to a system of algebraic equations that can be solved to obtain the unknown parameters β , λ and μ . Hence, the solution considered in Eq.(5) is obtained [1].

3 Application

Example 1 Gardner equation

$$u_t - 6(u + \epsilon^2 u^2)u_x + u_{xxx} = 0. \quad (7)$$

By using the wave variable $\xi = x - ct$ and $u(x, t) = f(\xi)$ then equation becomes

$$-\frac{df(\xi)}{d\xi} - 6(f(\xi) + \epsilon^2 f^2(\xi)) \frac{df(\xi)}{d\xi} + \frac{d^3f(\xi)}{d\xi^3} = 0, \quad (8)$$

Intergrating Eq.(8)gives

$$-cf(\xi) - 6\frac{f^2(\xi)}{2} - 6\epsilon^2\frac{f^3(\xi)}{3} + \frac{d^2f(\xi)}{d\xi^2} = 0. \quad (9)$$

Substituting Eq.(6) into Eq.(9) gives:

$$\begin{aligned} & -\lambda \cos^\beta(\mu\xi) - 3\lambda^2 \cos^{2\beta}(\mu\xi) - 2\epsilon^2\lambda^3 \cos^{3\beta}(\mu\xi) - \lambda\beta\mu^2 \cos^\beta(\mu\xi) \\ & + \lambda\beta(\beta - 1)\mu^2 \cos^\beta(\mu\xi) - \lambda\beta(\beta - 1)\mu^2 \cos^\beta(\mu\xi) = 0, \end{aligned} \quad (10)$$

By equating the exponents and the coefficients of each pair of the cosine function we obtain the following system of algebraic equations:

$$\begin{cases} -c\lambda - \lambda\beta\mu^2 - \lambda\beta(\beta - 1)\mu^2 = 0, \\ -2\epsilon^2\lambda^3 + \lambda\beta(\beta - 1)\mu^2 = 0. \end{cases} \quad (11)$$

$$3\beta = \beta - 2 \rightarrow \beta = 1.$$

By using maple for solving the system Eq(11):

$$\beta = 1-, \mu = \pm i\sqrt{c}, \lambda = \pm \frac{\sqrt{c}}{\epsilon} \quad (12)$$

Substituting Eq.(12) into Eq.(5) gives: (figure1)

$$\begin{aligned} u(x, t) &= \pm \frac{\sqrt{c}^{-1}}{\cos} (\pm i\sqrt{c}(x - ct)), \\ u(x, t) &= \pm \frac{\sqrt{c}}{\cosh} (\sqrt{c}(x - ct)). \end{aligned} \quad (13)$$

Example 2

(2+1)-dimensional breaking soliton system

$$\begin{aligned} u_t + 4buv_x + 4bu_xv + bv_{xxy} &= 0, \\ v_x - u_y &= 0. \end{aligned} \quad (14)$$

Suppose $u(x, y, t) = f(\xi)$ and $v(x, y, t) = g(\xi)$ and $\xi = x + ky - ct$ then

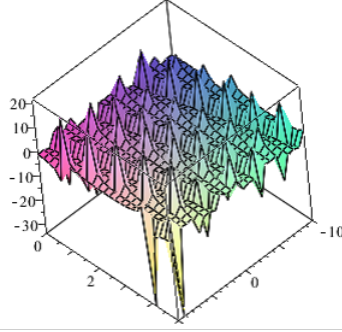


Fig. 1. shows the soliton solution for Gardner equation with increase time.
the (2+1)-dimensional soliton system becomes

$$-c \frac{df(\xi)}{d\xi} + 4bf(\xi) \frac{dg(\xi)}{d\xi} + 4b \frac{df(\xi)}{d\xi} g(\xi) + bk \frac{d^3 f(\xi)}{d\xi^3} = 0, \quad (15)$$

$$\frac{dg(\xi)}{d\xi} - k \frac{df(\xi)}{d\xi} = 0, \quad (16)$$

Integrating Eq.(15), Eq.(16) gives:

$$-cf(\xi) + 4bf(\xi)g(\xi) + bk \frac{d^2 f(\xi)}{d\xi^2} = 0, \quad (17)$$

$$g(\xi) - kf(\xi) = 0, \quad (18)$$

From Eq.(18)

$$g(\xi) = kf(\xi) = 0, \quad (19)$$

Substituting Eq.(19) into Eq.(17) gives:

$$-cf(\xi) + 4bkf^2(\xi) + bk \frac{d^2 f(\xi)}{d\xi^2} = 0. \quad (20)$$

Substituting Eq.(6) into Eq.(20) gives:

$$\begin{aligned} & -c\lambda \cos^\beta(\mu\xi) + 4bk\lambda^2 \cos^{2\beta}(\mu\xi) + bk[-\lambda\beta\mu^2 \cos^\beta(\mu\xi) \\ & + \lambda\mu^2\beta(\beta-1) \cos^{\beta-2}(\mu\xi) \\ & - \lambda\mu^2\beta(\beta-1) \cos^\beta(\mu\xi)] = 0, \end{aligned} \quad (21)$$

By equating the exponents and the coefficients of each pair of the cosine function we obtain the following system of algebraic equations:

$$\begin{cases} -c\lambda - bk\lambda\mu^2\beta - bk\lambda\mu^2\beta(\beta - 1) = 0, \\ 4bk\lambda^2 + bk\lambda\mu^2\beta(\beta - 1) = 0. \end{cases} \quad (22)$$

$$2\beta = \beta - 2 \rightarrow \beta = 2.$$

By using Maple for solving the system Eq.(22) we get:

$$\beta = -2, \quad \mu = \pm \frac{i}{2} \sqrt{\frac{c}{bk}}, \quad \lambda = \frac{3}{8} \frac{c}{bk} \quad (23)$$

Then by substituting Eq.(23) into Eq.(5) then, the exact soliton solutions of the (2+1)-dimensional breaking soliton system can be written in the form: (figure 2, 3)

$$u(x, y, t) = \frac{3}{8} \frac{c}{bk} \cos^{-2} \left(\pm \frac{i}{2} \sqrt{\frac{c}{bk}} (x + ky - ct) \right),$$

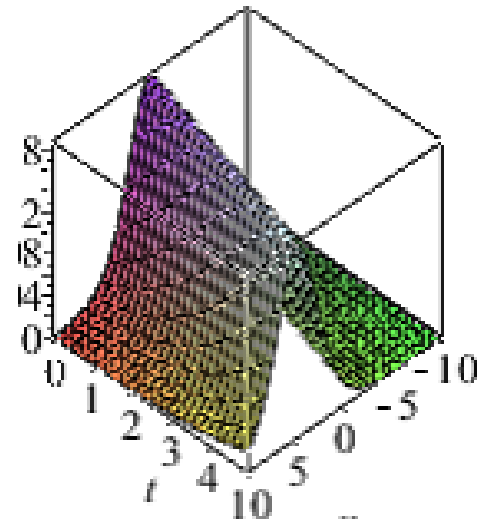
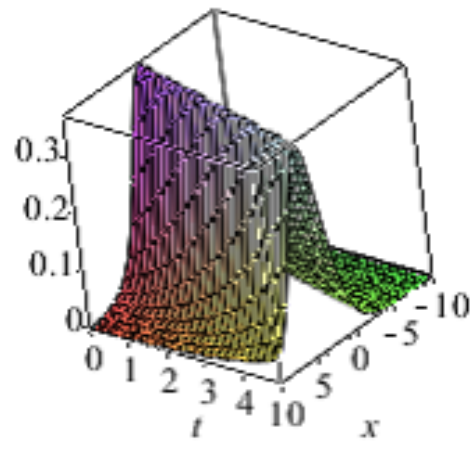
$$u(x, y, t) = \frac{3}{8} \frac{c}{bk} \sec^2 h^2 \left(\frac{1}{2} \sqrt{\frac{c}{bk}} (x + ky - ct) \right). \quad (24)$$

$$v(x, y, t) = \frac{3}{8} \frac{c}{b} \sec^2 h^2 \left(\frac{1}{2} \sqrt{\frac{c}{bk}} (x + ky - ct) \right). \quad (25)$$

Figures 3 and 3 show soliton solutions u, v of the (2+1)-dimensional breaking soliton system at $y = 0$ to $b = 1, k = 2$

4 Conclusion

In the letter, the cosine function method has been successfully applied to find the solution for two nonlinear partial differential equations such as Gardner and (2+1)-dimensional breaking soliton system. The cosine function method is used to find a new exact solution.



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